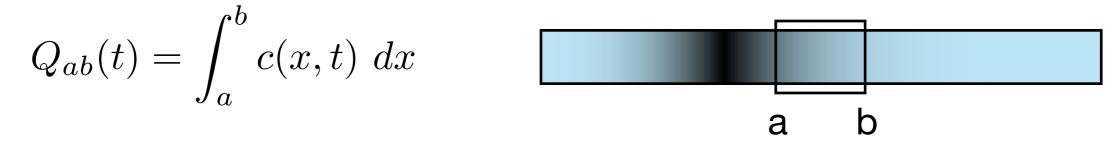
Today

- Diffusion equation -
 - derivation (transport eqns in general)
 - initial conditions, boundary conditions
 - steady state
 - separation of variables

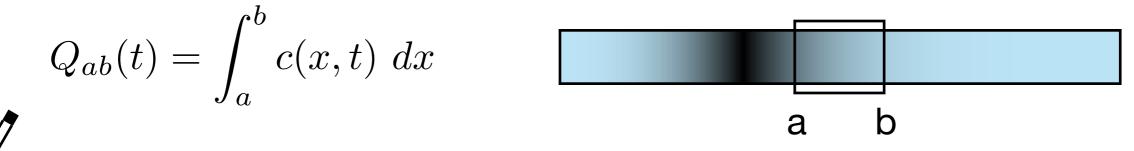
Conservation equations

c(x,t) is linear mass density of ink in a long narrow tube.



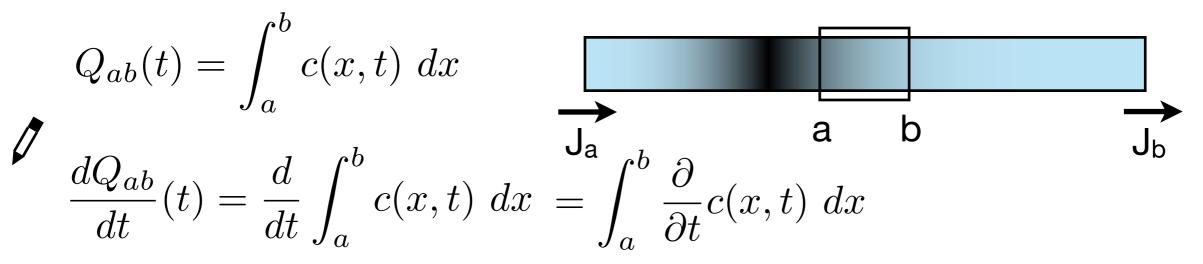
Conservation equations

c(x,t) is linear mass density of ink in a long narrow tube.



Conservation equations

c(x,t) is linear mass density of ink in a long narrow tube.

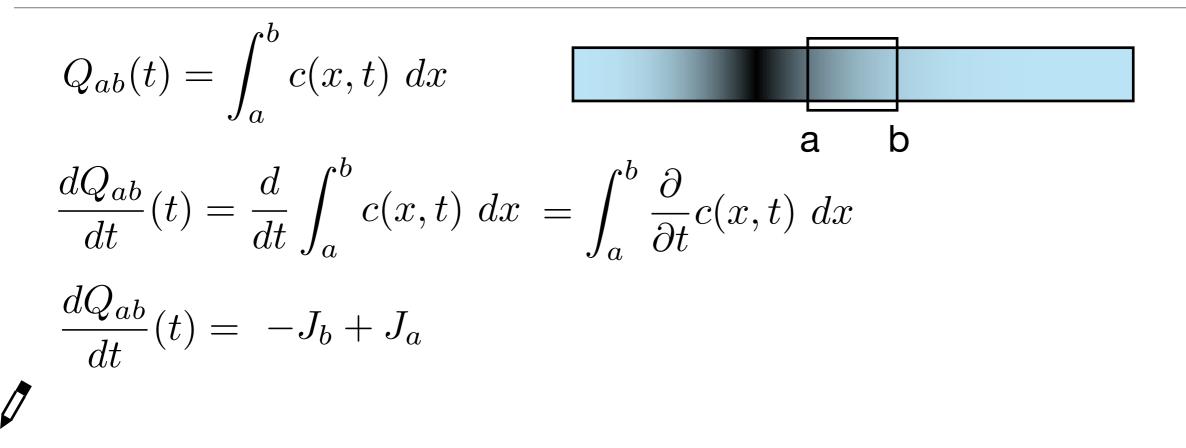


Define the flux J_a to be the amount of mass crossing the line x=a per unit of time (particles moving right count as positive flux).

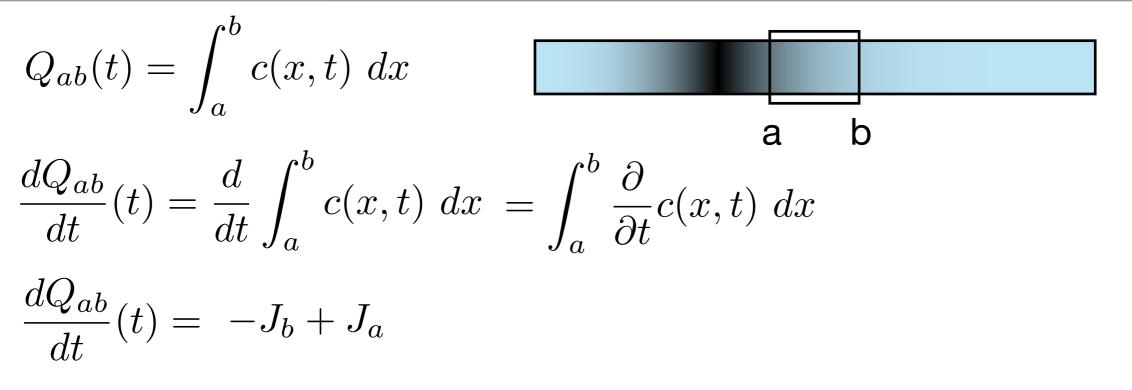
In that case, the change of Q inside the a-b box can also be counted watching flux, that is, flux at a - flux at b:

$$\frac{dQ_{ab}}{dt}(t) = -J_b + J_a$$

Conservation equations - Transport equation



Conservation equations - Transport equation



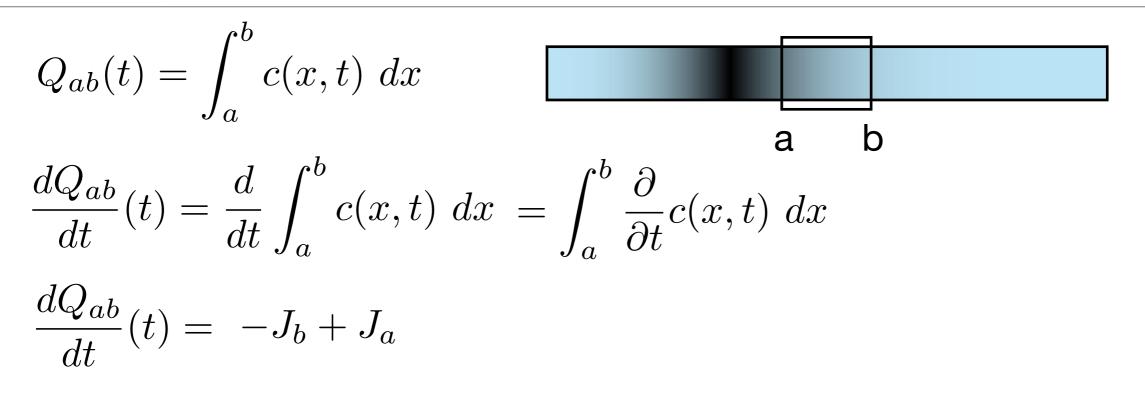
Need a model for flux. Let's consider simpler case first (not diffusion yet!) If fluid in pipe is moving with velocity v, flux is vc: $J_a = vc(a, t)$

$$\frac{dQ_{ab}}{dt}(t) = -J_b + J_a = -vc(b,t) + vc(a,t) = -vc(x,t) \Big|_a^b = -\int_a^b v \frac{\partial c}{\partial x} dx$$

$$\int_{a}^{b} \frac{\partial}{\partial t} c(x,t) \, dx = -\int_{a}^{b} v \frac{\partial c}{\partial x} \, dx \implies \frac{\partial c}{\partial t} = -v \frac{\partial c}{\partial x}$$

Called Transport equation.

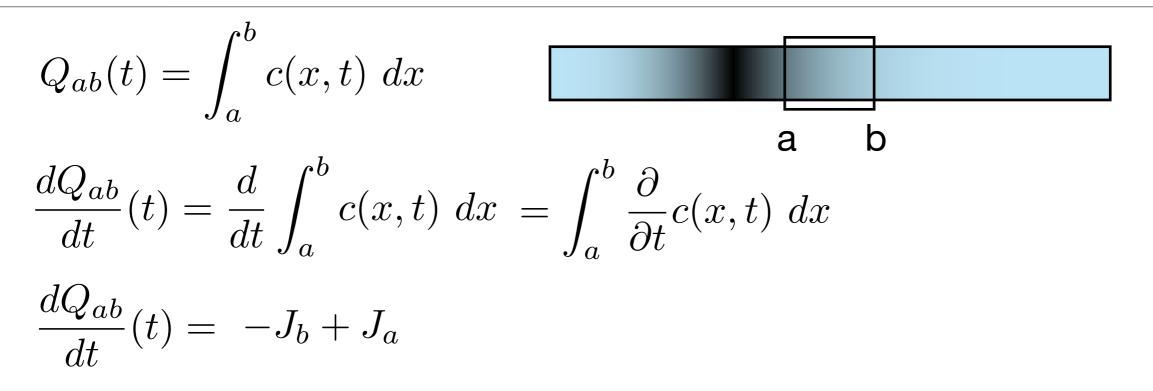
Conservation equations - Diffusion equation



Now lets consider diffusion. We can derive this from chemical potential but it also makes sense that for diffusion: $J_a = -D \left. \frac{\partial c}{\partial x} \right|_{x=0}$

$$b$$
 induce concertation and other $b_a - D$

Conservation equations - Diffusion equation



Now lets consider diffusion. We can derive this from chemical potential but it also makes sense that for diffusion: $J_a = -D \left. \frac{\partial c}{\partial x} \right|_{x=a}$ $\int \frac{dQ_{ab}}{dQ_{ab}}(t) = -J_b + J_a = D \left. \frac{\partial c}{\partial x} \right|_{x=a} - D \left. \frac{\partial c}{\partial x} \right|_{x=a} = D \left. \frac{\partial c}{\partial x} \right|_{x=a}$

$$\frac{\partial t}{\partial x}\Big|_{x=b} \qquad \frac{\partial x}{\partial x}\Big|_{x=a} \qquad \frac{\partial x}{\partial x}\Big|_{a}$$
$$\int_{a}^{b} \frac{\partial}{\partial t}c(x,t) \, dx = \int_{a}^{b} D \frac{\partial^{2}c}{\partial x^{2}} \, dx \qquad \Rightarrow \quad \frac{\partial}{\partial t}c(x,t) = D \frac{\partial^{2}}{\partial x^{2}}c(x,t)$$

Conservation equations - Diffusion equation

$$Q_{ab}(t) = \int_{a}^{b} c(x,t) dx$$

$$\frac{dQ_{ab}}{dt}(t) = \frac{d}{dt} \int_{a}^{b} c(x,t) dx = \int_{a}^{b} \frac{\partial}{\partial c} c(x,t) dx$$

$$\frac{dQ_{ab}}{dt}(t) = -J_{b} + J_{a}$$
The Diffusion Equation
$$\frac{dc}{dt} = D \frac{d^{2}c}{dx^{2}}$$
rom chemical potential but
it also makes sense that for diffusion:
$$J_{a} = -D \frac{\partial c}{\partial x}\Big|_{x=a}$$

$$\frac{dQ_{ab}}{dt}(t) = -J_{b} + J_{a} = D \frac{\partial c}{\partial x}\Big|_{x=b} - D \frac{\partial c}{\partial x}\Big|_{x=a} = D \frac{\partial c}{\partial x}\Big|_{a}$$

$$\int_{a}^{b} \frac{\partial}{\partial t}c(x,t) dx = \int_{a}^{b} D \frac{\partial^{2}c}{\partial x^{2}} dx \quad \Rightarrow \quad \frac{\partial}{\partial t}c(x,t) = D \frac{\partial^{2}}{\partial x^{2}}c(x,t)$$

• One derivative in time requires an initial condition in t.

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.
- Boundary conditions:

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.
- Boundary conditions:
 - $c(0,t) = c_0$ and $c(L,t) = c_L$ Dirichlet conditions

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.
- Boundary conditions:
 - $c(0,t) = c_0$ and $c(L,t) = c_L$ Dirichlet conditions • $\frac{dc}{dx}(0,t) = m_0$ and $\frac{dc}{dx}(L,t) = m_L$ Neumann conditions

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.
- Boundary conditions:
 - $c(0,t) = c_0$ and $c(L,t) = c_L$ **Dirichlet conditions** • $\frac{dc}{dr}(0,t) = m_0$ and $\frac{dc}{dr}(L,t) = m_L$ Neumann conditions Neumann conditions also Recall flux: $J_a = -D \frac{dc}{dr}(a,t)$ called flux conditions (no-

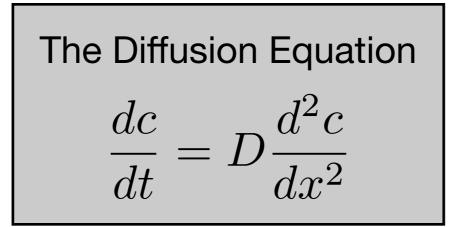
flux when $m_0 = m_L = 0$)

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.
- Boundary conditions:

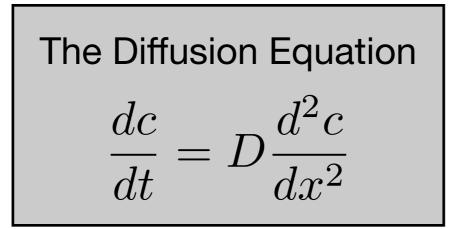
•
$$c(0,t) = c_0$$
 and $c(L,t) = c_L$
• $\frac{dc}{dx}(0,t) = m_0$ and $\frac{dc}{dx}(L,t) = m_L$
• $c(0,t) = c_0$ and $\frac{dc}{dx}(L,t) = m_L$
• $c(0,t) = c_0$ and $\frac{dc}{dx}(L,t) = m_L$
Mixed conditions

- One derivative in time requires an initial condition in t.
- Two derivatives in space require two "initial conditions" in x (i.e. one at x=0 and one at x=L). Called boundary conditions (BCs).
- Initial condition: c(x,0) = f(x) where f(x) gives initial concentration profile.
- Boundary conditions:

•
$$c(0,t) = c_0$$
 and $c(L,t) = c_L$
• $\frac{dc}{dx}(0,t) = m_0$ and $\frac{dc}{dx}(L,t) = m_L$
• $c(0,t) = c_0$ and $\frac{dc}{dx}(L,t) = m_L$
• $a\frac{dc}{dx}(0,t) + bc(0,t) = m_0$
Birichlet conditions
Neumann conditions
(no-flux conditions)
Mixed conditions
Robin conditions

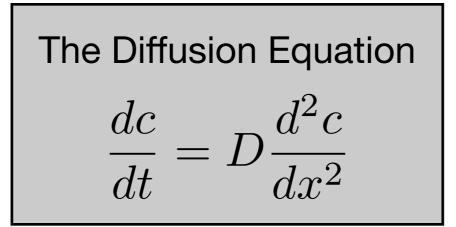


What does a steady state of the Diffusion equation look like?



What does a steady state of the Diffusion equation look like?

$$0 = D \frac{d^2c}{dx^2}$$



What does a steady state of the Diffusion equation look like?

$$0 = D \frac{d^2 c}{dx^2}$$
$$c_{ss}(x) = Ax + B$$

The Diffusion Equation $\frac{dc}{dt} = D \frac{d^2c}{dx^2}$

What does a steady state of the Diffusion equation look like?

$$0 = D \frac{d^2 c}{dx^2}$$
$$c_{ss}(x) = Ax + B$$

 A and B can be determined using the BCs. Getting A from Neumann conditions requires using the IC as well (total mass conservation).

Separation of variables

Separation of variables

• Doc cam

	Deriving the FS coefficient formulae
	C (1 1 1 1 1 1 for periodic functions (with period P)
	Define the dot product for periodic interiors (control f $f(x) \circ g(x) = \int f(x) \cdot g(x) dx = \int_{period}^{p_{x}} f(x) g(x) dx$
	Let $V_n(x) = cos(2\pi nx)$, $W_n(x) = sin(2\pi nx)$, $V_o(x) = 1$. $(n = 1, 2,)$
This pdf is also posted	Recall (or calculate for yourself) that $V_0(x) \circ V_0(x) = P$, $V_m(x) \circ V_n(x) = 0$ for $m \neq n$, $V_n(x) \circ V_n(x) = P_2$ $W_m(x) \circ V_n(x) = 0$, $W_m(x) \circ W_n(x) = 0$ for $m \neq n$, $W_n(x) \circ W_n(x) = P_2$
on the lecture slides page.	Suppose $f(x)$ can be represented exactly as a FS. Thus $f(x) = A_0 V_0(x) + \sum_{m=1}^{\infty} a_m V_m(x) + \sum_{m=1}^{\infty} b_m V_m(x).$
	Find its FS coefficients. As with vectors use "o" to find Ao, an, bn.
	To find A_{o_1} $f(x)_{o_1} V_{o_2}(x) = A_{o_1} V_{o_2}(x) + \sum_{n=1}^{\infty} a_n V_m(x)_{o_1} V_{o_2}(x) + \sum_{n=1}^{\infty} b_n W_m(x)_{o_1} V_{o_2}(x) = A_{o_1} P$
	$f(x) \circ V_{0}(x) = A_{0}V_{0}(x) \circ V_{0}(x) + \sum_{\substack{M=1 \\ M=1}}^{\infty} a_{M}V_{M}(x) \circ V_{0}(x) + \sum_{\substack{M=1 \\ M=1}}^{\infty} b_{M}W_{M}(x) \circ V_{0}(x) = A_{0} \cdot P$ Thus, $A_{0} = \frac{1}{p} f(x) \circ V_{0}(x) = \frac{1}{p} \int_{-p_{1}}^{p_{2}} f(x) dx$.
	To find an, $f(x) \circ V_n(x) = A_N(x) \circ V_n(x) + \sum_{n=1}^{\infty} Q_n V_n(x) \circ V_n(x) + \sum_{n=1}^{\infty} Q_n V_n(x) \circ V_n(x) = Q_n V_n(x) \circ V_n(x)$
	$f(x) \circ V_{n}(x) = A_{\sigma} V_{\sigma}(x) + \sum_{m=1}^{\infty} a_{m} V_{m}(x) \circ V_{n}(x) + \sum_{m=1}^{\infty} b_{m} W_{m}(x) \circ V_{n}(x) = a_{n} V_{n}(x) \circ V_{n}(x)$ $Thus, a_{n} = \frac{2}{p} f(x) \circ V_{n}(x) = \frac{2}{p} \int_{-f_{2}}^{f_{2}} f(x) \cos \frac{2\pi i T x}{p} dx.$
	Similarly, $b_n = \frac{2}{p} \int_{-\frac{1}{2}}^{\frac{1}{2}} f(x) \le \ln \frac{2n\pi x}{p} dx$
	In many cases, we will have P=2L (but not always!) So
	$A_{b} = \frac{1}{2L} \int_{-L}^{L} f(x) dx$
	$\alpha_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) dx$
	$b_{n} = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) dx$

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = A_0 + a_1 \cos\left(\frac{\pi x}{L}\right) + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = A_0 + a_1 \cos\left(\frac{\pi x}{L}\right) + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

メ

-1

 Our hope is that f(x) = f_{FS}(x) so we calculate coefficients as if they were equal:

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = A_0 + a_1 \cos\left(\frac{\pi x}{L}\right) + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

_

-1

 Our hope is that f(x) = f_{FS}(x) so we calculate coefficients as if they were equal:

$$A_0 = \frac{1}{2L} \int_{-L}^{L} f(x) \, dx$$
$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) \, dx$$
$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) \, dx$$

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = A_0 + a_1 \cos\left(\frac{\pi x}{L}\right) + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

 Our hope is that f(x) = f_{FS}(x) so we calculate coefficients as if they were equal:

$$A_0 = \frac{1}{2L} \int_{-L}^{L} f(x) \, dx$$
$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) \, dx$$
$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) \, dx$$

• To simplify formulas, usually define

$$a_0 = 2A_0 = \frac{1}{L} \int_{-L}^{L} f(x) \, dx$$

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = A_0 + a_1 \cos\left(\frac{\pi x}{L}\right) + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

 Our hope is that f(x) = f_{FS}(x) so we calculate coefficients as if they were equal:

$$A_0 = \frac{1}{2L} \int_{-L}^{L} f(x) \, dx$$
$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) \, dx$$
$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) \, dx$$

• To simplify formulas, usually define

-1

$$a_0 = 2A_0 = \frac{1}{L} \int_{-L}^{L} f(x) \, dx$$

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = \underbrace{\frac{a_0}{2}}_{+b_1 \sin\left(\frac{\pi x}{L}\right)} + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

 Our hope is that f(x) = f_{FS}(x) so we calculate coefficients as if they were equal:

$$A_0 = \frac{1}{2L} \int_{-L}^{L} f(x) \, dx$$
$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) \, dx$$
$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) \, dx$$

• To simplify formulas, usually define

-1

$$a_0 = 2A_0 = \frac{1}{L} \int_{-L}^{L} f(x) \, dx$$

• Find the Fourier series for $f(x) = 2u_0(x)-1$ on the interval [-1,1].

$$f_{FS}(x) = \underbrace{\frac{a_0}{2}}_{+b_1 \sin\left(\frac{\pi x}{L}\right)} + a_2 \cos\left(\frac{2\pi x}{L}\right) + \cdots$$
$$+b_1 \sin\left(\frac{\pi x}{L}\right) + b_2 \sin\left(\frac{2\pi x}{L}\right) + \cdots$$

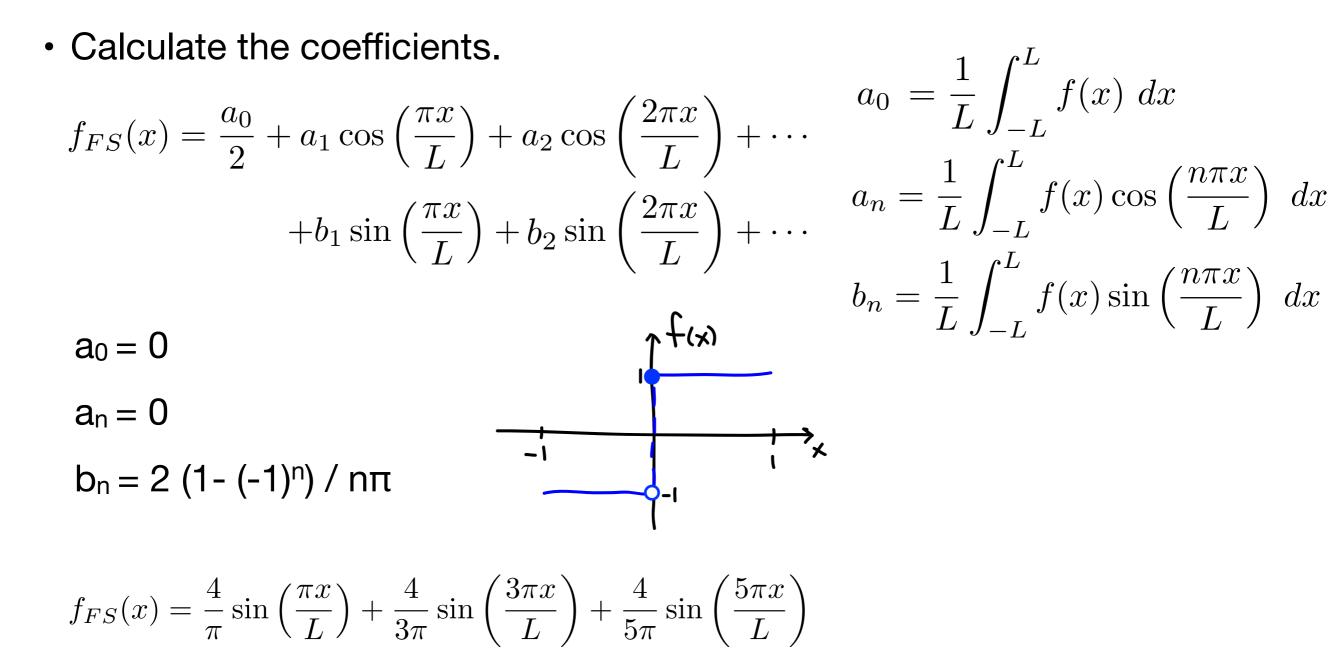
 Our hope is that f(x) = f_{FS}(x) so we calculate coefficients as if they were equal:

$$A_{0} = \frac{1}{2L} \int_{-L}^{L} f(x) \, dx \quad \begin{array}{l} A_{0} \text{ is the average} \\ \text{value of } f(x)! \end{array}$$
$$a_{n} = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) \, dx$$
$$b_{n} = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) \, dx$$

• To simplify formulas, usually define

-1

$$a_0 = 2A_0 = \frac{1}{L} \int_{-L}^{L} f(x) \, dx$$



https://www.desmos.com/calculator/tlvtikmi0y

Does $f(x) = f_{FS}(x)$ for all x? Problems at jumps! x=-1, 0, 1

• **Theorem** Suppose f and f' are piecewise continuous on [-L,L] and periodic beyond that interval. Then $f(x) = f_{FS}(x)$ at all points at which f is continuous. Furthermore, at points of discontinuity, $f_{FS}(x)$ takes the value of the midpoint of the jump. That is,

$$f_{FS}(x) = \frac{f(x^+) + f(x^-)}{2}$$

Heat/Diffusion equation - example

Find the solution to the heat/diffusion equation

$$u_t = 7u_{xx}$$

subject to BCs

$$u(0,t) = 0 = u(4,t)$$

and with IC

$$u(x,0) = \begin{cases} 1, & 0 \le x \le 2\\ 0, & 2 < x \le 4 \end{cases}$$

• The "warming up the milk bottle" example.

https://www.desmos.com/calculator/zvowjmu30g

 $u_t = 4u_{xx}$

$$u(0,t) = u(2,t) = 0$$
$$u(x,0) = x$$

https://www.desmos.com/calculator/yt7kztckeu

https://www.desmos.com/calculator/wcdvgrveez

 $u_t = 4u_{xx}$

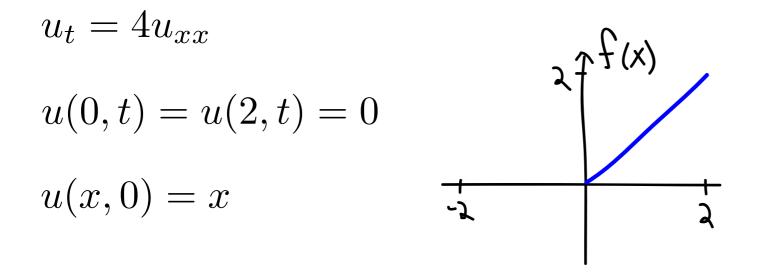
$$u(0,t) = u(2,t) = 0$$
$$u(x,0) = x$$

(A)
$$u(x,t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n e^{-n^2 \pi^2 t} \cos \frac{n \pi x}{2}$$
 $a_0 = 1, \ a_n = -\frac{8}{n^2 \pi^2} \text{ for } n \text{ even}$
(0 for $n \text{ odd}$)
(B) $u(x,t) = \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 t} \sin \frac{n \pi x}{2}$ $b_n = \frac{(-1)^{n+1} 4}{n \pi}$

 $u_t = 4u_{xx}$

$$u(0,t) = u(2,t) = 0$$
$$u(x,0) = x$$

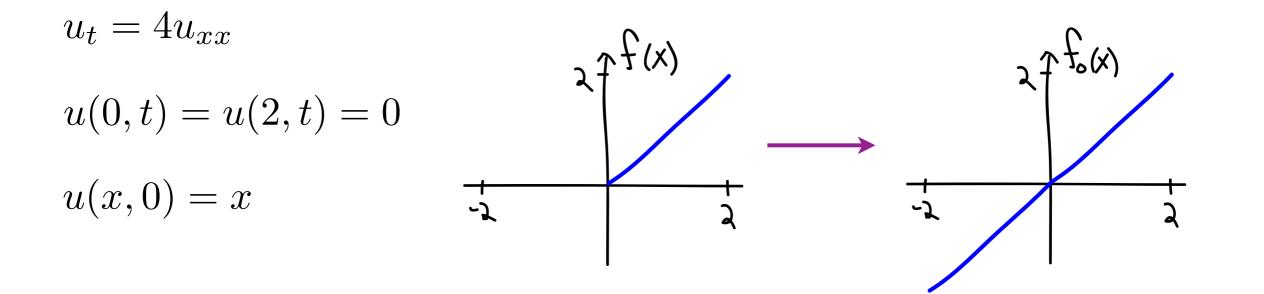
(A)
$$u(x,t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n e^{-n^2 \pi^2 t} \cos \frac{n \pi x}{2}$$
 $a_0 = 1, \ a_n = -\frac{8}{n^2 \pi^2} \text{ for } n \text{ even}$
(0 for $n \text{ odd}$)
(C) for $n \text{ odd}$
(B) $u(x,t) = \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 t} \sin \frac{n \pi x}{2}$ $b_n = \frac{(-1)^{n+1} 4}{n \pi}$



(A)
$$u(x,t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n e^{-n^2 \pi^2 t} \cos \frac{n \pi x}{2}$$
 $a_0 = 1, \ a_n = -\frac{8}{n^2 \pi^2} \text{ for } n \text{ even}$
(0 for $n \text{ odd}$)

$$\bigstar(\mathsf{B}) \ u(x,t) = \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 t} \sin \frac{n\pi x}{2}$$

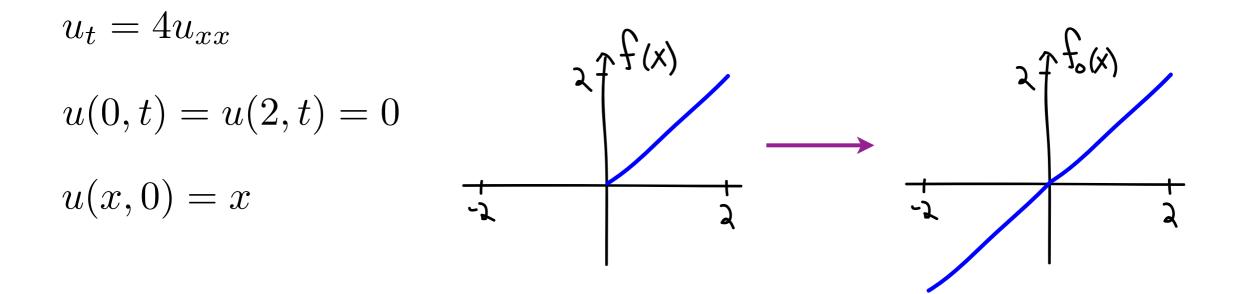
 $b_n = \frac{\sqrt{n\pi}}{n\pi}$



(A)
$$u(x,t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n e^{-n^2 \pi^2 t} \cos \frac{n \pi x}{2}$$
 $a_0 = 1, \ a_n = -\frac{8}{n^2 \pi^2}$ for n even (0 for n odd)

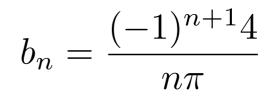
$$\bigstar(\mathsf{B}) \ u(x,t) = \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 t} \sin \frac{n\pi x}{2}$$

 $b_n = \frac{(-1)^{n+1}4}{n\pi}$



(A)
$$u(x,t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n e^{-n^2 \pi^2 t} \cos \frac{n \pi x}{2}$$
 $a_0 = 1, \ a_n = -\frac{8}{n^2 \pi^2}$ for n even
(0 for n odd)

$$\bigstar(\mathsf{B}) \ u(x,t) = \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 t} \sin \frac{n \pi x}{2}$$



Show Desmos movies.

https://www.desmos.com/calculator/yt7kztckeu

https://www.desmos.com/calculator/wcdvgrveez

$$\begin{aligned} u_t &= 4u_{xx} \\ \frac{du}{dx} \Big|_{x=0,2} &= 0 \\ u(x,0) &= \cos \frac{3\pi x}{2} \end{aligned}$$

 $\begin{aligned} u_t &= 4u_{xx} \\ \frac{du}{dx} \Big|_{x=0,2} &= 0 \\ u(x,0) &= \cos\frac{3\pi x}{2} \end{aligned}$

The IC is an eigenvector! Note that it satisfies the BCs.

 $\begin{aligned} u_t &= 4u_{xx} \\ \frac{du}{dx} \Big|_{x=0,2} &= 0 \\ u(x,0) &= \cos\frac{3\pi x}{2} \end{aligned}$

The IC is an eigenvector! Note that it satisfies the BCs.
$$v_3(x) = \cos \frac{3\pi x}{2}$$

The IC is an eigenvector! Note that it satisfies the BCs.

 $n\pi x$

$$\frac{du}{dx}\Big|_{x=0,2} = 0 \qquad \qquad v_3(x) = \cos\frac{3\pi x}{2}$$
$$u(x,0) = \cos\frac{3\pi x}{2} \qquad \qquad v_n(x) = \cos\frac{n\pi x}{2}$$

 $u_t = 4u_{xx}$

 $\begin{array}{l} u_t = 4u_{xx} \\ \left. \frac{du}{dx} \right|_{x=0,2} = 0 \\ u(x,0) = \cos \frac{3\pi x}{2} \end{array} \end{array}$ The IC is an eigenvector! Note that it satisfies the BCs. $v_3(x) = \cos \frac{3\pi x}{2} \\ v_n(x) = \cos \frac{\pi x}{2} \\ u_n(x,t) = e^{\lambda_n t} \cos \frac{n\pi x}{2} \end{array}$

 $u_t = 4u_{xx}$ The IC is an eigenvector! Note that it satisfies the BCs. $\left. \frac{du}{dx} \right|_{x=0,2} = 0$ $u(x,0) = \cos\frac{3\pi x}{2}$

$$v_{3}(x) = \cos \frac{3\pi x}{2}$$
$$v_{n}(x) = \cos \frac{n\pi x}{2}$$
$$u_{n}(x,t) = e^{\lambda_{n}t} \cos \frac{n\pi x}{2}$$
$$\frac{\partial}{\partial t}u_{n}(x,t) = \lambda_{n}e^{\lambda_{n}t} \cos \frac{n\pi x}{2}$$

The IC is an eigenvector! Note that it satisfies the BCs. $\left.\frac{du}{dx}\right|_{x=0,2} = 0$ $u(x,0) = \cos\frac{3\pi x}{2}$

 $u_t = 4u_{xx}$

$$v_{3}(x) = \cos \frac{3\pi x}{2}$$

$$v_{n}(x) = \cos \frac{n\pi x}{2}$$

$$u_{n}(x,t) = e^{\lambda_{n}t} \cos \frac{n\pi x}{2}$$

$$\frac{\partial}{\partial t}u_{n}(x,t) = \lambda_{n}e^{\lambda_{n}t} \cos \frac{n\pi x}{2}$$

$$\frac{\partial^{2}}{\partial x^{2}}u_{n}(x,t) = -\frac{n^{2}\pi^{2}}{4}e^{\lambda_{n}t} \cos \frac{n\pi x}{2}$$

 $u_t = 4u_{xx}$ The IC is an eigenvector! Note that it satisfies the BCs. $\left. \frac{du}{dx} \right|_{x=0,2} = 0$ $v_3(x) = \cos\frac{3\pi x}{2}$ $u(x,0) = \cos\frac{3\pi x}{2}$ $v_n(x) = \cos\frac{n\pi x}{2}$ $u_n(x,t) = e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $\frac{\partial}{\partial t}u_n(x,t) = \lambda_n e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $4\frac{\partial^2}{\partial x^2}u_n(x,t) = -\frac{4n^2\pi^2}{4}e^{\lambda_n t}\cos\frac{n\pi x}{2}$

 $u_t = 4u_{xx}$ The IC is an eigenvector! Note that it satisfies the BCs. $\left. \frac{du}{dx} \right|_{x=0.2} = 0$ $v_3(x) = \cos\frac{3\pi x}{2}$ $u(x,0) = \cos\frac{3\pi x}{2}$ $v_n(x) = \cos\frac{n\pi x}{2}$ $u_n(x,t) = e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $\frac{\partial}{\partial t} u_n(x,t) = \lambda_n e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $4\frac{\partial^2}{\partial r^2}u_n(x,t) = \frac{An^2\pi^2}{A}e^{\lambda_n t}\cos\frac{n\pi x}{2}$

 $u_t = 4u_{xx}$ The IC is an eigenvector! Note that it satisfies the BCs. $\left. \frac{du}{dx} \right|_{x=0,2} = 0$ $v_3(x) = \cos\frac{3\pi x}{2}$ $u(x,0) = \cos\frac{3\pi x}{2}$ $v_n(x) = \cos\frac{n\pi x}{2}$ $u_n(x,t) = e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $\frac{\partial}{\partial t}u_n(x,t) = \lambda_n e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $4\frac{\partial^2}{\partial x^2}u_n(x,t) = \frac{An^2\pi^2}{N}e^{\lambda_n t}\cos\frac{n\pi x}{2}$

 $u_t = 4u_{xx}$ The IC is an eigenvector! Note that it satisfies the BCs. $\left. \frac{du}{dx} \right|_{x=0,2} = 0$ $v_3(x) = \cos\frac{3\pi x}{2}$ $u(x,0) = \cos\frac{3\pi x}{2}$ $v_n(x) = \cos\frac{n\pi x}{2}$ $u_n(x,t) = e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $\frac{\partial}{\partial t}u_n(x,t) = \lambda_n e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $4\frac{\partial^2}{\partial x^2}u_n(x,t) = \frac{An^2\pi^2}{A}e^{\lambda_n t}\cos\frac{n\pi x}{2}$ $\lambda_n = -n^2 \pi^2$

 $u_t = 4u_{xx}$ The IC is an eigenvector! Note that it satisfies the BCs. $\left. \frac{du}{dx} \right|_{x=0,2} = 0$ $v_3(x) = \cos\frac{3\pi x}{2}$ $u(x,0) = \cos\frac{3\pi x}{2}$ $v_n(x) = \cos\frac{n\pi x}{2}$ $u_n(x,t) = e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $\frac{\partial}{\partial t}u_n(x,t) = \lambda_n e^{\lambda_n t} \cos \frac{n\pi x}{2}$ $4\frac{\partial^2}{\partial x^2}u_n(x,t) = \frac{An^2\pi^2}{A}e^{\lambda_n t}\cos\frac{n\pi x}{2}$ $\lambda_n = -n^2 \pi^2$ So the solution is $u(x,t) = e^{-9\pi^2 t} \cos \frac{3\pi x}{2}$