Today

- Review of solutions to the Diffusion Equation with various BCs.
- The Wave Equation.
- Separation of variables.

 $u_{t} = Du_{xx}$ u(0,t) = u(L,t) = 0 u(x,0) = f(x) f(x) f(x) u(x,0) = f(x)







$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L} + \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}$$



• All coefficients will be non-zero. Not particularly useful for solving the BCs.













$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos \frac{n\pi x}{L} dx$$



$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos \frac{n\pi x}{L} \, dx$$



$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos \frac{n\pi x}{L} \, dx$$



$$a_n = 0$$



$$a_n = 0 \qquad \qquad b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin \frac{n\pi x}{L} dx$$



$$a_n = 0$$
 $b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin \frac{n\pi x}{L} dx$



$$a_n = 0 \qquad \qquad b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} \, dx$$



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 $u_t = Du_{xx}$

u(0,t) = u(L,t) = 0u(x,0) = f(x)

$$u(x,t) = \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 D t/L^2} \sin \frac{n\pi x}{L}$$
$$b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

 $u_t = Du_{xx}$

$$\left. \frac{\partial u}{\partial x} \right|_{x=0,L} = 0$$

u(x,0) = f(x)

$$u(x,t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n e^{-n^2 \pi^2 D t/L^2} \cos \frac{n\pi x}{L}$$

$$a_n = \frac{2}{L} \int_0^L f(x) \cos \frac{n\pi x}{L} \, dx$$

 $u_t = Du_{xx}$ u(0,t) = au(L,t) = bu(x,0) = f(x)

$$u(x,t) = a + \frac{b-a}{L}x + \sum_{n=1}^{\infty} b_n e^{-n^2 \pi^2 D t/L^2} \sin \frac{n\pi x}{L}$$
$$b_n = \frac{2}{L} \int_0^L \left(f(x) - a - \frac{b-a}{L}x \right) \sin \frac{n\pi x}{L} \, dx$$

 Adding the linear function to the usual solution to the Dirichlet problem ensures that the BCs are satisfied without changing the fact that it satisfies the PDE.

$$u_t = Du_{xx}$$
$$u(0,t) = 0$$
$$\frac{\partial u}{\partial x}\Big|_{x=L} = 0$$
$$u(x,0) = f(x)$$







 $u_t = Du_{xx}$ u(0,t) = 0 $\frac{\partial u}{\partial x}\Big|_{x=L} = 0$ u(x,0) = f(x)



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Review of solutions to the Diffusion Equation



Review of solutions to the Diffusion Equation



Motion of a string

• Motion of a string

Motion of a string



• Motion of a string

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Motion of a string



m = mass of string,

L = length,

 $\rho = m/L =$ "density",

T = tension in string,

 $u_n = vertical displacement of nth mass.$

Motion of a string



m = mass of string, L = length,

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Motion of a string



 $\frac{m}{3}\frac{\partial^2 u_2}{\partial t^2} = \frac{T}{\Delta x}((u_3 - u_2) - (u_2 - u_1))$

m = mass of string, L = length, ρ = m/L = "density",

T = tension in string, $u_n = vertical displacement$ of nth mass.

(not obvious - requires more detail)

Motion of a string



 $\begin{array}{ll} m = mass \ of \ string, & T = \\ L = length, & u_n = \\ \rho = m/L = "density", & of n \end{array}$

T = tension in string, u_n = vertical displacement of nth mass.

 $\frac{m}{n}\frac{\partial^2 u_k}{\partial t^2} = \frac{T}{\Delta x}((u_{k+1} - u_k) - (u_k - u_{k-1})) \quad \text{(not obvious - requires more detail)}$

Motion of a string



T = tension in string, $u_n =$ vertical displacement of nth mass.

 $\frac{m}{n}\frac{\partial^2 u_k}{\partial t^2} = \frac{T}{\Delta x}((u_{k+1} - u_k) - (u_k - u_{k-1})) \text{ (not obvious - requires more detail)}$ $n = L/\Delta x - 1 \approx L/\Delta x$

Motion of a string



 $\begin{array}{ll} m = mass \ of \ string, \\ L = length, \\ \rho = m/L = \ "density", \end{array} \begin{array}{ll} T = tension \ in \ string, \\ u_n = vertical \ displacement \\ of \ n^{th} \ mass. \end{array}$

 $\frac{m}{n}\frac{\partial^2 u_k}{\partial t^2} = \frac{T}{\Delta x}((u_{k+1} - u_k) - (u_k - u_{k-1})) \text{ (not obvious - requires more detail)} \\ \frac{m\Delta x}{L}\frac{\partial^2 u_k}{\partial t^2} = \frac{T}{\Delta x}((u_{k+1} - u_k) - (u_k - u_{k-1}))$

Motion of a string



 $\begin{array}{ll} m = mass \ of \ string, \\ L = length, \\ \rho = m/L = \ "density", \end{array} \begin{array}{ll} T = tension \ in \ string, \\ u_n = vertical \ displacement \\ of \ n^{th} \ mass. \end{array}$

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Motion of a string



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$$\frac{\partial^2 u_k}{\partial t^2} = \frac{T}{\rho} \frac{\frac{u_{k+1} - u_k}{\Delta x} - \frac{u_k - u_{k-1}}{\Delta x}}{\Delta x}$$

Motion of a string



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$$\frac{\partial^2 u_k}{\partial t^2} = \frac{1}{\rho} \frac{\Delta x}{\Delta x} \xrightarrow{\Delta x} \Delta x \qquad \longrightarrow \qquad \frac{\partial^2 u}{\partial t^2} = \frac{1}{\rho} \frac{\partial^2 u}{\partial x^2}$$

Motion of a string



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$$\frac{\partial^2 u_k}{\partial t^2} = \frac{1}{\rho} \frac{\overline{\Delta x}}{\Delta x} \frac{\overline{\Delta x}}{\Delta x} \qquad \xrightarrow{\Delta x \to 0} \qquad \frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

Motion of a string



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Motion of a string



 $\begin{array}{ll} m = mass \ of \ string, & T = tension \ in \ string, \\ L = length, & u_n = vertical \ displacement \\ \rho = m/L = "density", & of \ n^{th} \ mass. \end{array}$

$$\begin{aligned} \frac{m}{n} \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\Delta x} ((u_{k+1} - u_k) - (u_k - u_{k-1})) & \text{(not obvious - requires more detail)} \\ n &= L/\Delta x - 1 \approx L/\Delta x \\ \frac{m\Delta x}{L} \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\Delta x} ((u_{k+1} - u_k) - (u_k - u_{k-1})) \\ \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\rho \Delta x^2} ((u_{k+1} - u_k) - (u_k - u_{k-1})) & k = 1, 2, 3, \dots, n \\ \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\rho} \frac{\frac{u_{k+1} - u_k}{\Delta x} - \frac{u_k - u_{k-1}}{\Delta x}}{\Delta x} \xrightarrow{\Delta x \to 0} & \frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} & \begin{array}{c} \text{BCs:} \\ u(0, t) &= 0 \\ u(L, t) &= 0 \end{array} \end{aligned}$$

Motion of a string



ICs:

 $\begin{array}{ll} m = mass \ of \ string, & T = tension \ in \ string, \\ L = length, & u_n = vertical \ displacement \\ \rho = m/L = "density", & of \ n^{th} \ mass. \end{array}$

$$\begin{aligned} \frac{m}{n} \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\Delta x} ((u_{k+1} - u_k) - (u_k - u_{k-1})) & \text{(not obvious - requires more detail)} \\ n &= L/\Delta x - 1 \approx L/\Delta x \\ \frac{m\Delta x}{L} \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\Delta x} ((u_{k+1} - u_k) - (u_k - u_{k-1})) \\ \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\rho \Delta x^2} ((u_{k+1} - u_k) - (u_k - u_{k-1})) & k = 1, 2, 3, \dots, n \\ \frac{\partial^2 u_k}{\partial t^2} &= \frac{T}{\rho} \frac{\frac{u_{k+1} - u_k}{\Delta x} - \frac{u_k - u_{k-1}}{\Delta x}}{\Delta x} \xrightarrow{\Delta x \to 0} & \frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} & \underset{u(0, t) = 0}{\text{BCs:}} \\ u(0, t) &= 0 \end{aligned}$$

Motion of a string



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- Other physical systems described by the wave equation:
- Sound waves u(x,t) is the air pressure, c is speed of sound.
- Water waves u(x,t) is water height, c is wave speed.
- Electromagnetic waves u(x,t) is field intensity, c is speed of light ...

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

• There is a way to calculate eigenvalues as we did for the Diffusion Equation but that requires rewriting it as a first order system:

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find e of t

find evalues/vectors of this operator

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fi

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• Inspired by the fundamental solutions to the Diffusion Equation, assume

 $u_n(x,t) = T(t)v(x)$

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Inspired by the fundamental solutions to the Diffusion Equation, assume

$$u_n(x,t) = T(t)v(x) \qquad \qquad u_n(x,t) = e^{\lambda_n t} \sin \frac{n\pi x}{L}$$

 There is a way to calculate eigenvalues as we did for the Diffusion Equation but that requires rewriting it as a first order system:

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 - $u_n(x,t) = T(t)v(x)$ $u_n(x,t) = e^{\lambda_n t} \sin \frac{n\pi x}{L}$ $T(t) = e^{\lambda t}$

 There is a way to calculate eigenvalues as we did for the Diffusion Equation but that requires rewriting it as a first order system:

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 - $$\begin{split} u_n(x,t) &= T(t)v(x) \\ u_n(x,t) &= e^{\lambda_n t} \sin \frac{n\pi x}{L} \\ T(t) &= e^{\lambda t} \\ \text{Expect oscillations so} \\ T(t) &= e^{i\beta t} \end{split}$$

(or sines and cosines)

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$$u_n(x,t) = T(t)v(x)$$
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$$T''(t)v(x) = c^2 T(t)v''(x)$$
$$\frac{T''(t)}{T(t)} = c^2 \frac{v''(x)}{v(x)}$$

$$u_n(x,t) = T(t)v(x)$$
$$\frac{\partial^2 u_n}{\partial t^2} = T''(t)v(x)$$
$$\frac{\partial^2 u_n}{\partial x^2} = T(t)v''(x)$$
$$T''(t)v(x) = c^2 T(t)v''(x)$$
$$\frac{T''(t)}{T(t)} = c^2 \frac{v''(x)}{v(x)} = \alpha$$

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$$\begin{split} u_n(x,t) &= T(t)v(x) \qquad \alpha > 0 \Rightarrow \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) \\ T''(t)v(x) &= c^2 T(t)v''(x) \\ \frac{T''(t)}{T(t)} &= c^2 \frac{v''(x)}{v(x)} = \alpha \\ T''(t) &= \alpha T(t) \\ v''(x) &= \frac{\alpha}{c^2}v(x) \end{split}$$

$$u_n(x,t) = T(t)v(x) \qquad \alpha > 0 \Rightarrow \quad T(t) = e^{\pm\sqrt{\alpha}t}$$

$$\frac{\partial^2 u_n}{\partial t^2} = T''(t)v(x)$$

$$\frac{\partial^2 u_n}{\partial x^2} = T(t)v''(x)$$

$$T''(t)v(x) = c^2T(t)v''(x)$$

$$\frac{T''(t)}{T(t)} = c^2\frac{v''(x)}{v(x)} = \alpha$$

$$T''(t) = \alpha T(t)$$

$$v''(x) = \frac{\alpha}{c^2}v(x)$$

$$u_{n}(x,t) = T(t)v(x) \qquad \alpha > 0 \Rightarrow T(t) = e^{\pm\sqrt{\alpha}t}$$

$$\frac{\partial^{2}u_{n}}{\partial t^{2}} = T''(t)v(x) \qquad v(x) = e^{\pm\frac{\sqrt{\alpha}}{c}x}$$

$$\frac{\partial^{2}u_{n}}{\partial x^{2}} = T(t)v''(x)$$

$$T''(t)v(x) = c^{2}T(t)v''(x)$$

$$\frac{T''(t)}{T(t)} = c^{2}\frac{v''(x)}{v(x)} = \alpha$$

$$T''(t) = \alpha T(t)$$

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$$u_n(x,t) = T(t)v(x)$$

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$$T''(t)v(x) = c^2 T(t)v''(x)$$

$$\frac{T''(t)}{T(t)} = c^2 \frac{v''(x)}{v(x)} = \alpha$$

$$T''(t) = \alpha T(t)$$

$$v''(x) = \frac{\alpha}{c^2}v(x)$$

$$\begin{array}{ll} \alpha > 0 \Rightarrow & T(t) = e^{\pm \sqrt{\alpha}t} \\ & v(x) = e^{\pm \frac{\sqrt{\alpha}}{c}x} \end{array}$$

- no oscillations in time,
- can't satisfy BCs! (v(0)=v(L)=0)

$$u_n(x,t) = T(t)v(x) \qquad \alpha > 0 \Rightarrow$$
$$\frac{\partial^2 u_n}{\partial t^2} = T''(t)v(x) \qquad \bullet$$
$$\frac{\partial^2 u_n}{\partial t^2} = T''(t)v(x) \qquad \bullet$$

$$\frac{\partial^{n} u_{n}}{\partial x^{2}} = T(t)v''(x) \qquad \qquad \alpha =$$

$$T''(t)v(x) = c^2 T(t)v''(x)$$

> 0
$$\Rightarrow$$
 $T(t) = e^{\pm \sqrt{\alpha}t}$
 $v(x) = e^{\pm \frac{\sqrt{\alpha}}{c}x}$

- no oscillations in time,
- can't satisfy BCs! (v(0)=v(L)=0)

$$= 0$$
 doesn't work either (check).

$$\frac{T''(t)}{T(t)} = c^2 \frac{v''(x)}{v(x)} = \alpha$$
$$T''(t) = \alpha T(t)$$
$$v''(x) = \frac{\alpha}{c^2} v(x)$$

 $T''(t) = \alpha T(t)$

Inspired by the fundamental solutions to the Diffusion Equation, assume

$$\begin{split} u_n(x,t) &= T(t)v(x) \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) \\ T''(t)v(x) &= c^2 T(t)v''(x) \\ \frac{T''(t)}{T(t)} &= c^2 \frac{v''(x)}{v(x)} &= \alpha \end{split} \qquad \alpha > 0 \Rightarrow T(t) = e^{\pm \sqrt{\alpha}t} \\ v(x) &= e^{\pm \sqrt{\alpha}c} x \\ \bullet \text{ no oscillations in time,} \\ \bullet \text{ can't satisfy BCs! (v(0)=v(L)=0)} \\ \alpha &= 0 \quad \text{ doesn't work either (check).} \\ \alpha < 0 \Rightarrow \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t + A_2 \sin \sqrt{\alpha} \\ \bullet x \\ T(t) &= A_1 \cos \sqrt{-\alpha}t$$

 $-\alpha t$

$$v''(x) = \frac{\alpha}{c^2}v(x)$$

$$\begin{split} u_n(x,t) &= T(t)v(x) & \alpha > 0 \Rightarrow \quad T(t) = e^{\pm\sqrt{\alpha}t} \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) & \nu(x) = e^{\pm\frac{\sqrt{\alpha}}{c}x} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ T''(t)v(x) &= c^2T(t)v''(x) & \alpha < 0 \Rightarrow \\ \frac{T''(t)}{T(t)} &= c^2\frac{v''(x)}{v(x)} &= \alpha & v(x) = B_1\cos\sqrt{-\alpha}t + A_2\sin\sqrt{-\alpha}t \\ \frac{T''(t)}{T(t)} &= \alpha T(t) & v''(x) = \frac{\alpha}{c^2}v(x) \end{split}$$

$$\begin{split} u_n(x,t) &= T(t)v(x) & \alpha > 0 \Rightarrow \quad T(t) = e^{\pm\sqrt{\alpha}t} \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) & \nu(x) = e^{\pm\frac{\sqrt{\alpha}}{c}x} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ T''(t)v(x) &= c^2T(t)v''(x) & \alpha < 0 \Rightarrow \\ \frac{T''(t)}{T(t)} &= c^2\frac{v''(x)}{v(x)} &= \alpha & v(x) = B_1\cos\sqrt{-\alpha}t + A_2\sin\sqrt{-\alpha}t \\ \frac{T''(t)}{T(t)} &= \alpha T(t) & u(0,t) = 0 \\ v''(x) &= \frac{\alpha}{c^2}v(x) & u(0,t) = 0 \end{split}$$

• Inspired by the fundamental solutions to the Diffusion Equation, assume

$$\begin{split} u_n(x,t) &= T(t)v(x) & \alpha > 0 \Rightarrow \quad T(t) = e^{\pm\sqrt{\alpha}t} \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) & \nu(x) = e^{\pm\frac{\sqrt{\alpha}}{c}x} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ T''(t)v(x) &= c^2T(t)v''(x) & \alpha < 0 \Rightarrow \\ \frac{T''(t)}{T(t)} &= c^2\frac{v''(x)}{v(x)} &= \alpha & v(x) = B_1\cos\sqrt{-\alpha}t + A_2\sin\sqrt{-\alpha}t \\ \frac{T''(t)}{T(t)} &= \alpha T(t) & u(0,t) = 0 \\ v''(x) &= \frac{\alpha}{c^2}v(x) & u(0,t) = 0 \end{split}$$

 ${\mathcal X}$

$$\begin{split} u_n(x,t) &= T(t)v(x) & \alpha > 0 \Rightarrow \quad T(t) = e^{\pm\sqrt{\alpha}t} \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) & \nu(x) = e^{\pm\frac{\sqrt{\alpha}}{c}x} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha < 0 \text{ solutions in time,} \\ T''(t)v(x) &= c^2T(t)v''(x) & \alpha < 0 \Rightarrow \\ \frac{T''(t)}{T(t)} &= c^2\frac{v''(x)}{v(x)} &= \alpha & v(x) = B_1\cos\sqrt{-\alpha}t + A_2\sin\sqrt{-\alpha}t \\ \frac{T''(t)}{T(t)} &= \alpha T(t) & u(0,t) = 0 \\ v''(x) &= \frac{\alpha}{c^2}v(x) & u(L,t) = 0 \end{split}$$

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$$\begin{split} u_n(x,t) &= T(t)v(x) & \alpha > 0 \Rightarrow T(t) = e^{\pm\sqrt{\alpha}t} \\ \frac{\partial^2 u_n}{\partial t^2} &= T''(t)v(x) & \nu(x) = e^{\pm\frac{\sqrt{\alpha}}{c}x} \\ \frac{\partial^2 u_n}{\partial x^2} &= T(t)v''(x) & \alpha = 0 & \text{doesn't work either (check).} \\ T''(t)v(x) &= c^2T(t)v''(x) & \alpha < 0 \Rightarrow \\ T''(t) &= c^2\frac{v''(x)}{v(x)} &= \alpha & v(x) = B_1\cos\sqrt{-\alpha}t + A_2\sin\sqrt{-\alpha}t \\ \frac{T''(t)}{T(t)} &= c^2\frac{v''(x)}{v(x)} &= \alpha & v(x) = B_1\cos\sqrt{-\alpha}x + B_2\sin\frac{\sqrt{-\alpha}}{c}x \\ T''(t) &= \alpha T(t) & u(0,t) = 0 \\ v''(x) &= \frac{\alpha}{c^2}v(x) & u_n(x,t) = \left(A_n\cos\frac{n\pi ct}{L} + B_n\sin\frac{n\pi ct}{L}\right)\sin\frac{n\pi x}{L} \end{split}$$

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$
$$u(0,t) = 0$$
$$u(L,t) = 0$$
$$u(x,0) = f(x)$$
$$\frac{\partial}{\partial t}u(x,0) = g(x)$$

https://www.desmos.com/calculator/dmfijt0e6q

 $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ u(0,t) = 0u(L,t) = 0u(x,0) = f(x) $\frac{\partial}{\partial t} u(x,0) = g(x)$

$$u_n(x,t) = \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L}\right) \sin \frac{n\pi x}{L}$$

https://www.desmos.com/calculator/dmfijt0e6q

 $\begin{aligned} \frac{\partial^2 u}{\partial t^2} &= c^2 \frac{\partial^2 u}{\partial x^2} \\ u(0,t) &= 0 \\ u(L,t) &= 0 \\ u(x,0) &= f(x) \end{aligned} \qquad u_n(x,t) = \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L}\right) \sin \frac{n\pi x}{L} \\ u(x,0) &= f(x) \\ \frac{\partial}{\partial t} u(x,0) &= g(x) \end{aligned}$

 $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ u(0,t) = 0u(L,t) = 0u(x,0) = f(x) $\frac{\partial}{\partial t}u(x,0) = g(x)$

$$u_n(x,t) = \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L}\right) \sin \frac{n\pi x}{L}$$
$$u(x,t) = \sum_{n=1}^{\infty} \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L}\right) \sin \frac{n\pi x}{L}$$

• Pull, hold and let go of a guitar string:

 $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ u(0,t) = 0u(L,t) = 0u(x,0) = f(x) $\frac{\partial}{\partial t}u(x,0) = g(x)$

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• Pull, hold and let go of a guitar string: $\frac{\partial}{\partial t}u(x,0) = 0 \implies B_n = 0$

 $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ u(0,t) = 0u(L,t) = 0u(x,0) = f(x) $\frac{\partial}{\partial t}u(x,0) = g(x)$

$$u_n(x,t) = \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L}\right) \sin \frac{n\pi x}{L}$$
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 $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ u(0,t) = 0u(L,t) = 0u(x,0) = f(x) $\frac{\partial}{\partial t}u(x,0) = g(x)$

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$$u(x,t) = \sum_{n=1}^{\infty} \left(A_n \cos \frac{n\pi ct}{L} + B_n \sin \frac{n\pi ct}{L}\right) \sin \frac{n\pi x}{L}$$

• Pull, hold and let go of a guitar string: $\frac{\partial}{\partial t}u(x,0) = 0 \quad \Rightarrow B_n = 0$ $u(x,t) = \sum_{n=1}^{\infty} A_n \cos \frac{n\pi ct}{L} \sin \frac{n\pi x}{L}$

$$u(x,t) = \sum_{n=1}^{\infty} b_n \left(\sin\left(\frac{n\pi}{L}(x-ct)\right) + \sin\left(\frac{n\pi}{L}(x+ct)\right) \right)$$

https://www.desmos.com/calculator/dmfijt0e6q